# Andrea Malchiodi (SISSA, Trieste)

New improved Moser-Trudinger inequalities and singular Liouville equations on compact surfaces

Geometric Properties for Parabolic and Elliptic PDEs Cortona, June 20-24, 2011

(1) A.M. - D.Ruiz: New improved Moser-Trudinger inequalities and singular Liouville equations on compact surfaces, GAFA, to appear.

(2) D.Bartolucci - F.De Marchis - A.M.: Supercritical conformal metrics on surfaces with conical singularities, IMRN, to appear.

(3) A.Carlotto - A.M.: Weighted barycentric sums and singular Liouville equations on compact surfaces, preprint, 2011.

We consider the following mean field equation on a compact orientable surface  $(\Sigma, g)$  without boundary:

...

$$(E_{\rho}) \qquad -\Delta u = \rho \left( k(x)e^{2u} - 1 \right) - 2\pi \sum_{j=1}^{m} \alpha_j \left( \delta_{p_j} - 1 \right).$$

We consider the following mean field equation on a compact orientable surface  $(\Sigma, g)$  without boundary:

$$(E_{\rho}) \qquad -\Delta u = \rho \left( k(x)e^{2u} - 1 \right) - 2\pi \sum_{j=1}^{m} \alpha_j \left( \delta_{p_j} - 1 \right).$$

Here  $\rho$  is a positive parameter,  $k : \Sigma \to \mathbb{R}$  a smooth positive function,  $\alpha_j \in \mathbb{R}$  and  $\delta_{p_j}$  is the Dirac mass at  $p_j \in \Sigma$ .

We consider the following mean field equation on a compact orientable surface  $(\Sigma, g)$  without boundary:

$$(E_{\rho}) \qquad -\Delta u = \rho \left( k(x)e^{2u} - 1 \right) - 2\pi \sum_{j=1}^{m} \alpha_j \left( \delta_{p_j} - 1 \right).$$

Here  $\rho$  is a positive parameter,  $k : \Sigma \to \mathbb{R}$  a smooth positive function,  $\alpha_j \in \mathbb{R}$  and  $\delta_{p_i}$  is the Dirac mass at  $p_j \in \Sigma$ .

• We assume also, without loss of generality, that  $|\Sigma| = 1$ .

In this model a wave function  $\psi:\mathbb{R}^2\to\mathbb{C}$  satisfying the NLS

$$i\psi_t = -\frac{1}{2m}\Delta\psi - \Gamma|\psi|^2\psi$$

is coupled to a Gauge field  $A_{\mu}$ 

In this model a wave function  $\psi:\mathbb{R}^2\to\mathbb{C}$  satisfying the NLS

$$i\psi_t = -\frac{1}{2m}\Delta\psi - \Gamma|\psi|^2\psi$$

is coupled to a Gauge field  $A_{\mu}$  ( $\partial_{\mu} \mapsto \partial_{\mu} - iA_{\mu}$ ).

In this model a wave function  $\psi:\mathbb{R}^2\to\mathbb{C}$  satisfying the NLS

$$i\psi_t = -\frac{1}{2m}\Delta\psi - \Gamma|\psi|^2\psi$$

is coupled to a Gauge field  $A_{\mu}$  ( $\partial_{\mu} \mapsto \partial_{\mu} - iA_{\mu}$ ). C-S coupling

$$F_{\mu,\nu} = \frac{1}{\kappa} \varepsilon_{\mu\nu\gamma} J^{\gamma}; \qquad F_{\mu\nu} = \partial_{\mu} A_{\nu} - \partial_{\nu} A_{\mu}.$$

In this model a wave function  $\psi:\mathbb{R}^2\to\mathbb{C}$  satisfying the NLS

$$i\psi_t = -\frac{1}{2m}\Delta\psi - \Gamma|\psi|^2\psi$$

is coupled to a Gauge field  $A_{\mu}$  ( $\partial_{\mu} \mapsto \partial_{\mu} - iA_{\mu}$ ). C-S coupling

$$F_{\mu,\nu} = \frac{1}{\kappa} \varepsilon_{\mu\nu\gamma} J^{\gamma}; \qquad F_{\mu\nu} = \partial_{\mu} A_{\nu} - \partial_{\nu} A_{\mu}.$$

Here  $\varepsilon_{\mu\nu\gamma}$  is the antisymmetric symbol and  $J^{\mu}$  is the *current* 

$$J^{\mu} = (\rho, \overrightarrow{J}); \qquad \rho = |\psi|^2, \quad \overrightarrow{J} = \frac{i}{2m} \left( \psi \nabla \overline{\psi} - \overline{\psi} \nabla \psi \right).$$

In this model a wave function  $\psi : \mathbb{R}^2 \to \mathbb{C}$  satisfying the NLS

$$i\psi_t = -\frac{1}{2m}\Delta\psi - \Gamma|\psi|^2\psi$$

is coupled to a Gauge field  $A_{\mu}$  ( $\partial_{\mu} \mapsto \partial_{\mu} - iA_{\mu}$ ). C-S coupling

$$F_{\mu,\nu} = \frac{1}{\kappa} \varepsilon_{\mu\nu\gamma} J^{\gamma}; \qquad F_{\mu\nu} = \partial_{\mu} A_{\nu} - \partial_{\nu} A_{\mu}.$$

Here  $\varepsilon_{\mu\nu\gamma}$  is the antisymmetric symbol and  $J^{\mu}$  is the *current* 

$$J^{\mu} = (\rho, \vec{J}); \qquad \rho = |\psi|^2, \quad \vec{J} = \frac{i}{2m} \left( \psi \nabla \overline{\psi} - \overline{\psi} \nabla \psi \right).$$

In the static self-dual regime,  $\Gamma = -\frac{1}{m\kappa}$ ,  $v = \log |\psi|$  satisfies

$$-\Delta v = \frac{1}{\kappa}e^{2v} - 2\pi \sum_{j=1}^{m} N_j \delta_{p_j},$$

where the  $p_j$ 's are the zeroes of  $\psi$  (vortices), with order  $N_j$ .

In this model a wave function  $\psi : \mathbb{R}^2 \to \mathbb{C}$  satisfying the NLS

$$i\psi_t = -\frac{1}{2m}\Delta\psi - \Gamma|\psi|^2\psi$$

is coupled to a Gauge field  $A_{\mu}$  ( $\partial_{\mu} \mapsto \partial_{\mu} - iA_{\mu}$ ). C-S coupling

$$F_{\mu,\nu} = \frac{1}{\kappa} \varepsilon_{\mu\nu\gamma} J^{\gamma}; \qquad F_{\mu\nu} = \partial_{\mu} A_{\nu} - \partial_{\nu} A_{\mu}.$$

Here  $\varepsilon_{\mu\nu\gamma}$  is the antisymmetric symbol and  $J^{\mu}$  is the *current* 

$$J^{\mu} = (\rho, \vec{J}); \qquad \rho = |\psi|^2, \quad \vec{J} = \frac{i}{2m} \left( \psi \nabla \overline{\psi} - \overline{\psi} \nabla \psi \right).$$

In the static self-dual regime,  $\Gamma = -\frac{1}{m\kappa}$ ,  $v = \log |\psi|$  satisfies

$$-\Delta v = \frac{1}{\kappa}e^{2v} - 2\pi \sum_{j=1}^{m} N_j \delta_{p_j},$$

where the  $p_j$ 's are the zeroes of  $\psi$  (vortices), with order  $N_j$ . Self-duality implies that vortices do not interact.

Consider a conformal metric  $\tilde{g}$  on  $\Sigma$ 

Consider a conformal metric  $\tilde{g}$  on  $\Sigma$ : if we set  $\tilde{g} = 2e^{2w}g$ , then the Gaussian curvature transforms according to the law

$$-\Delta_g w + K_g = K_{\tilde{g}} e^{2w}.$$

Consider a conformal metric  $\tilde{g}$  on  $\Sigma$ : if we set  $\tilde{g} = 2e^{2w}g$ , then the Gaussian curvature transforms according to the law

$$-\Delta_g w + K_g = K_{\tilde{g}} e^{2w}.$$

A Dirac delta on the right-hand side gives a singular metric with conical structure.



Consider a conformal metric  $\tilde{g}$  on  $\Sigma$ : if we set  $\tilde{g} = 2e^{2w}g$ , then the Gaussian curvature transforms according to the law

$$-\Delta_g w + K_g = K_{\tilde{g}} e^{2w}.$$

A Dirac delta on the right-hand side gives a singular metric with conical structure.



For a cone with opening angle  $\theta = 2\pi(1 + \alpha)$ , the Gaussian curvature on the sides is zero, and it is a Dirac mass or order  $-2\pi\alpha$  at the vertex.

Consider a conformal metric  $\tilde{g}$  on  $\Sigma$ : if we set  $\tilde{g} = 2e^{2w}g$ , then the Gaussian curvature transforms according to the law

$$-\Delta_g w + K_g = K_{\tilde{g}} e^{2w}.$$

A Dirac delta on the right-hand side gives a singular metric with conical structure.



For a cone with opening angle  $\theta = 2\pi(1 + \alpha)$ , the Gaussian curvature on the sides is zero, and it is a Dirac mass or order  $-2\pi\alpha$  at the vertex. For standard cones  $\alpha \in (-1, 0)$ .

From now on suppose for simplicity that there is only one singularity p with weight  $\alpha > 0$ .

From now on suppose for simplicity that there is only one singularity p with weight  $\alpha > 0$ .

Let  $\boldsymbol{w}$  be a solution of

(1) 
$$-\Delta w = \delta_p - 1.$$

From now on suppose for simplicity that there is only one singularity p with weight  $\alpha > 0$ .

Let  $\boldsymbol{w}$  be a solution of

(1) 
$$-\Delta w = \delta_p - 1.$$

Clearly,  $w(x) \sim -\frac{1}{2\pi} \log |x-p|$  for x close to p.

From now on suppose for simplicity that there is only one singularity p with weight  $\alpha > 0$ .

Let  $\boldsymbol{w}$  be a solution of

(1) 
$$-\Delta w = \delta_p - 1.$$

Clearly,  $w(x) \sim -\frac{1}{2\pi} \log |x - p|$  for x close to p. Using the substitution  $u \mapsto u + 2\pi \alpha w$ , we obtain the equivalent problem:

$$(\tilde{E}_{\rho}) \qquad -\Delta u = \rho \left( \frac{h(x)e^{2u}}{\int_{\Sigma} h(x)e^{2u}} - 1 \right); \qquad h(x) \sim dist(x,p)^{2\alpha}.$$

From now on suppose for simplicity that there is only one singularity p with weight  $\alpha > 0$ .

Let  $\boldsymbol{w}$  be a solution of

(1) 
$$-\Delta w = \delta_p - 1.$$

Clearly,  $w(x) \sim -\frac{1}{2\pi} \log |x - p|$  for x close to p. Using the substitution  $u \mapsto u + 2\pi \alpha w$ , we obtain the equivalent problem:

$$(\tilde{E}_{\rho}) \qquad -\Delta u = \rho \left( \frac{h(x)e^{2u}}{\int_{\Sigma} h(x)e^{2u}} - 1 \right); \qquad h(x) \sim dist(x,p)^{2\alpha}.$$

 $(\tilde{E}_{\rho})$  is the Euler-Lagrange equation for  $I_{\rho}: H^1(\Sigma) \to \mathbb{R}$  def. as

(2) 
$$I_{\rho}(u) = \int_{\Sigma} |\nabla u|^2 + 2\rho \int_{\Sigma} u - \rho \log \int_{\Sigma} h(x) e^{2u}.$$

Recall the classical Moser-Trudinger inequality

$$(\mathsf{M}-\mathsf{T}) \qquad \log \int_{\Sigma} e^{2(u-\overline{u})} \leq \frac{1}{4\pi} \int_{\Sigma} |\nabla u|^2 + C; \qquad u \in H^1(\Sigma).$$

Recall the classical Moser-Trudinger inequality

(M-T) 
$$\log \int_{\Sigma} e^{2(u-\overline{u})} \leq \frac{1}{4\pi} \int_{\Sigma} |\nabla u|^2 + C; \quad u \in H^1(\Sigma).$$

With a weight one has *Troyanov's inequality* ([Troyanov, '91]) (T)  $\log \int_{\Sigma} dist(x,p)^{2\alpha} e^{2(u-\overline{u})} \leq \frac{1}{4\pi \min\{1,1+\alpha\}} \int_{\Sigma} |\nabla u|^2 + C.$ 

Recall the classical Moser-Trudinger inequality

(M-T) 
$$\log \int_{\Sigma} e^{2(u-\overline{u})} \leq \frac{1}{4\pi} \int_{\Sigma} |\nabla u|^2 + C; \quad u \in H^1(\Sigma).$$

With a weight one has *Troyanov's inequality* ([Troyanov, '91]) (T)  $\log \int_{\Sigma} dist(x,p)^{2\alpha} e^{2(u-\overline{u})} \leq \frac{1}{4\pi \min\{1,1+\alpha\}} \int_{\Sigma} |\nabla u|^2 + C.$ (compared to Miyamoto's work here the energy is free)

Recall the classical Moser-Trudinger inequality

(M-T) 
$$\log \int_{\Sigma} e^{2(u-\overline{u})} \leq \frac{1}{4\pi} \int_{\Sigma} |\nabla u|^2 + C; \quad u \in H^1(\Sigma).$$

With a weight one has *Troyanov's inequality* ([Troyanov, '91]) (T)  $\log \int_{\Sigma} dist(x,p)^{2\alpha} e^{2(u-\overline{u})} \leq \frac{1}{4\pi \min\{1,1+\alpha\}} \int_{\Sigma} |\nabla u|^2 + C.$ (compared to Miyamoto's work here the energy is free)

**Consequence** ([Troyanov, '91]) For  $\rho < 4\pi \min\{1, 1 + \alpha\}$  the functional  $I_{\rho}$  is coercive and solutions can be found as global minima using the direct methods of Calculus of Variations.

Recall the classical Moser-Trudinger inequality

(M-T) 
$$\log \int_{\Sigma} e^{2(u-\overline{u})} \leq \frac{1}{4\pi} \int_{\Sigma} |\nabla u|^2 + C; \quad u \in H^1(\Sigma).$$

With a weight one has *Troyanov's inequality* ([Troyanov, '91]) (T)  $\log \int_{\Sigma} dist(x,p)^{2\alpha} e^{2(u-\overline{u})} \leq \frac{1}{4\pi \min\{1,1+\alpha\}} \int_{\Sigma} |\nabla u|^2 + C.$ (compared to Miyamoto's work here the energy is free)

**Consequence** ([Troyanov, '91]) For  $\rho < 4\pi \min\{1, 1 + \alpha\}$  the functional  $I_{\rho}$  is coercive and solutions can be found as global minima using the direct methods of Calculus of Variations.

Note that for  $\alpha < 0$  one has a worse constant since  $dist(x, p)^{2\alpha}$  is a singular function.

Recall the classical Moser-Trudinger inequality

(M-T) 
$$\log \int_{\Sigma} e^{2(u-\overline{u})} \leq \frac{1}{4\pi} \int_{\Sigma} |\nabla u|^2 + C; \quad u \in H^1(\Sigma).$$

With a weight one has *Troyanov's inequality* ([Troyanov, '91]) (T)  $\log \int_{\Sigma} dist(x,p)^{2\alpha} e^{2(u-\overline{u})} \leq \frac{1}{4\pi \min\{1,1+\alpha\}} \int_{\Sigma} |\nabla u|^2 + C.$ (compared to Miyamoto's work here the energy is free)

**Consequence** ([Troyanov, '91]) For  $\rho < 4\pi \min\{1, 1 + \alpha\}$  the functional  $I_{\rho}$  is coercive and solutions can be found as global minima using the direct methods of Calculus of Variations.

Note that for  $\alpha < 0$  one has a worse constant since  $dist(x,p)^{2\alpha}$  is a singular function. For  $\alpha > 0$  instead the sharp constant is the same as that for  $h \equiv 1$ .



Using standard bubbles of the form

$$\varphi_{\lambda,x}(y) := \log \frac{\lambda}{1 + \lambda^2 dist(x,y)^2}; \qquad x \neq p,$$

it is easy to see that  $\inf I_{\rho} = -\infty$ .

Using standard bubbles of the form

$$\varphi_{\lambda,x}(y) := \log \frac{\lambda}{1 + \lambda^2 dist(x,y)^2}; \qquad x \neq p,$$

it is easy to see that  $\inf I_{\rho} = -\infty$ . Not much was known about existence of solutions in this case.

Using standard bubbles of the form

$$\varphi_{\lambda,x}(y) := \log \frac{\lambda}{1 + \lambda^2 dist(x,y)^2}; \qquad x \neq p,$$

it is easy to see that  $\inf I_{\rho} = -\infty$ . Not much was known about existence of solutions in this case.

**Perturbative results** ([Del Pino-Esposito=Musso, '05]): multipeak solutions in a blow-up regime.

Using standard bubbles of the form

$$\varphi_{\lambda,x}(y) := \log \frac{\lambda}{1 + \lambda^2 dist(x,y)^2}; \qquad x \neq p,$$

it is easy to see that  $\inf I_{\rho} = -\infty$ . Not much was known about existence of solutions in this case.

**Perturbative results** ([Del Pino-Esposito=Musso, '05]): multipeak solutions in a blow-up regime. For regular case works by Del Pino, Kowalczyk, Musso, Ruf, ...
#### The case $\rho > 4\pi$

Using standard bubbles of the form

$$\varphi_{\lambda,x}(y) := \log \frac{\lambda}{1 + \lambda^2 dist(x,y)^2}; \qquad x \neq p,$$

it is easy to see that  $\inf I_{\rho} = -\infty$ . Not much was known about existence of solutions in this case.

**Perturbative results** ([Del Pino-Esposito=Musso, '05]): multipeak solutions in a blow-up regime. For regular case works by Del Pino, Kowalczyk, Musso, Ruf, ...

**On-going computation of the L-S degree** ([C.C.Chen - C.S. Lin, a first paper in '10]): via a refined blow-up analysis and finite dimensional reductions.

#### The case $\rho > 4\pi$

Using standard bubbles of the form

$$\varphi_{\lambda,x}(y) := \log \frac{\lambda}{1 + \lambda^2 dist(x,y)^2}; \qquad x \neq p,$$

it is easy to see that  $\inf I_{\rho} = -\infty$ . Not much was known about existence of solutions in this case.

**Perturbative results** ([Del Pino-Esposito=Musso, '05]): multipeak solutions in a blow-up regime. For regular case works by Del Pino, Kowalczyk, Musso, Ruf, ...

**On-going computation of the L-S degree** ([C.C.Chen - C.S. Lin, a first paper in '10]): via a refined blow-up analysis and finite dimensional reductions.

Our goal is to develop a global variational strategy to find general critical points of saddle type.

#### The case $\rho > 4\pi$

Using standard bubbles of the form

$$\varphi_{\lambda,x}(y) := \log \frac{\lambda}{1 + \lambda^2 dist(x,y)^2}; \qquad x \neq p,$$

it is easy to see that  $\inf I_{\rho} = -\infty$ . Not much was known about existence of solutions in this case.

**Perturbative results** ([Del Pino-Esposito=Musso, '05]): multipeak solutions in a blow-up regime. For regular case works by Del Pino, Kowalczyk, Musso, Ruf, ...

**On-going computation of the L-S degree** ([C.C.Chen - C.S. Lin, a first paper in '10]): via a refined blow-up analysis and finite dimensional reductions.

Our goal is to develop a global variational strategy to find general critical points of <u>saddle type</u>. This could be more direct and more general than L-S theory, as degree cancelations may occur.

Let  $\mathcal H$  be an Hilbert space, and  $f:\mathcal H\to\mathbb R$  a smooth functional

Let  $\mathcal{H}$  be an Hilbert space, and  $f : \mathcal{H} \to \mathbb{R}$  a smooth functional A general method to finding critical points of saddle type is to look at the topological properties of the sublevels.

Let  $\mathcal{H}$  be an Hilbert space, and  $f : \mathcal{H} \to \mathbb{R}$  a smooth functional A general method to finding critical points of saddle type is to look at the topological properties of the sublevels. A change in topology might suggest the presence of a critical point



Let  $\mathcal{H}$  be an Hilbert space, and  $f : \mathcal{H} \to \mathbb{R}$  a smooth functional A general method to finding critical points of saddle type is to look at the topological properties of the sublevels. A change in topology might suggest the presence of a critical point



Naively, if a < b and  $f : \mathcal{H} \to \mathbb{R}$  has no critical points in  $\{a \leq f \leq b\}$ , then using the gradient flow  $\{f \leq b\}$  can be smoothly deformed into  $\{f \leq a\}$ , keeping  $\{f \leq a\}$  fixed

Let  $\mathcal{H}$  be an Hilbert space, and  $f : \mathcal{H} \to \mathbb{R}$  a smooth functional A general method to finding critical points of saddle type is to look at the topological properties of the sublevels. A change in topology might suggest the presence of a critical point



Naively, if a < b and  $f : \mathcal{H} \to \mathbb{R}$  has no critical points in  $\{a \leq f \leq b\}$ , then using the gradient flow  $\{f \leq b\}$  can be smoothly deformed into  $\{f \leq a\}$ , keeping  $\{f \leq a\}$  fixed (deformation lemma).

Let  $\mathcal{H}$  be an Hilbert space, and  $f : \mathcal{H} \to \mathbb{R}$  a smooth functional A general method to finding critical points of saddle type is to look at the topological properties of the sublevels. A change in topology might suggest the presence of a critical point



Naively, if a < b and  $f : \mathcal{H} \to \mathbb{R}$  has no critical points in  $\{a \leq f \leq b\}$ , then using the gradient flow  $\{f \leq b\}$  can be smoothly deformed into  $\{f \leq a\}$ , keeping  $\{f \leq a\}$  fixed (deformation lemma).

#### Consequence:

Let  $\mathcal{H}$  be an Hilbert space, and  $f: \mathcal{H} \to \mathbb{R}$  a smooth functional A general method to finding critical points of saddle type is to look at the topological properties of the sublevels. A change in topology might suggest the presence of a critical point



Naively, if a < b and  $f : \mathcal{H} \to \mathbb{R}$  has no critical points in  $\{a \leq f \leq b\}$ , then using the gradient flow  $\{f \leq b\}$  can be smoothly deformed into  $\{f \leq a\}$ , keeping  $\{f \leq a\}$  fixed (deformation lemma).

**Consequence**: if one level does not deform into another, there should be a critical point between the two.

• We must be careful though: the presence of *asymptotes* might prevent deforming sublevels without critical points in between



• We must be careful though: the presence of *asymptotes* might prevent deforming sublevels without critical points in between



Some compactness criterion is needed (Palais-Smale condition or similar ones)

• Compactness along flow lines is somehow related to compactness of solutions to  $(\tilde{E}_{\rho})$ 

• Compactness along flow lines is somehow related to compactness of solutions to  $(\tilde{E}_{\rho})$  (want to avoid *critical points at infinity*)

• Compactness along flow lines is somehow related to compactness of solutions to  $(\tilde{E}_{\rho})$  (want to avoid *critical points at infinity*)

A standard tool in the field is <u>blow-up analysis</u>: rescale solutions in order to obtain *standard profiles*.

• Compactness along flow lines is somehow related to compactness of solutions to  $(\tilde{E}_{\rho})$  (want to avoid *critical points at infinity*)

A standard tool in the field is <u>blow-up analysis</u>: rescale solutions in order to obtain *standard profiles*. Here they are of two kinds



• Compactness along flow lines is somehow related to compactness of solutions to  $(\tilde{E}_{\rho})$  (want to avoid *critical points at infinity*)

A standard tool in the field is <u>blow-up analysis</u>: rescale solutions in order to obtain *standard profiles*. Here they are of two kinds



**Theorem** ([Bartolucci-Tarantello, '02]) Suppose  $u_n$  are solutions of  $(\tilde{E}_{\rho_n})$ .

• Compactness along flow lines is somehow related to compactness of solutions to  $(\tilde{E}_{\rho})$  (want to avoid *critical points at infinity*)

A standard tool in the field is <u>blow-up analysis</u>: rescale solutions in order to obtain *standard profiles*. Here they are of two kinds



**Theorem** ([Bartolucci-Tarantello, '02]) Suppose  $u_n$  are solutions of  $(\tilde{E}_{\rho_n})$ . Then either  $u_n$  stays bounded in  $C^2(\Sigma)$  or it blows up k spheres,  $k \ge 0$ , [plus possibly an AF at p].

• Compactness along flow lines is somehow related to compactness of solutions to  $(\tilde{E}_{\rho})$  (want to avoid *critical points at infinity*)

A standard tool in the field is <u>blow-up analysis</u>: rescale solutions in order to obtain *standard profiles*. Here they are of two kinds



**Theorem** ([Bartolucci-Tarantello, '02]) Suppose  $u_n$  are solutions of  $(\tilde{E}_{\rho_n})$ . Then either  $u_n$  stays bounded in  $C^2(\Sigma)$  or it blows up k spheres,  $k \ge 0$ , [plus possibly an AF at p]. If blow-up occurs then one has  $\rho_n \to \overline{\rho} = 4k\pi [+4(1 + \alpha)\pi]$ .

 $\Lambda := 4\pi \mathbb{N} \cup (4\pi \mathbb{N} + 4\pi(1 + \alpha)).$ 

$$\Lambda := 4\pi\mathbb{N} \cup (4\pi\mathbb{N} + 4\pi(1+\alpha)).$$

**Consequence** (compactness): if  $\rho \notin \Lambda$ , then we have that

$$\Lambda := 4\pi \mathbb{N} \cup (4\pi \mathbb{N} + 4\pi(1+\alpha)).$$

**Consequence** (compactness): if  $\rho \notin \Lambda$ , then we have that

(1) solutions of  $(\tilde{E}_{\rho})$  stay bounded

$$\Lambda := 4\pi \mathbb{N} \cup (4\pi \mathbb{N} + 4\pi(1+\alpha)).$$

**Consequence** (compactness): if  $\rho \notin \Lambda$ , then we have that

(1) solutions of  $(\tilde{E}_{\rho})$  stay bounded (and have bounded energy)

$$\Lambda := 4\pi \mathbb{N} \cup (4\pi \mathbb{N} + 4\pi(1+\alpha)).$$

**Consequence** (compactness): if  $\rho \notin \Lambda$ , then we have that

(1) solutions of  $(\tilde{E}_{\rho})$  stay bounded (and have bounded energy)

(2) the deformation lemma holds true

$$\Lambda := 4\pi \mathbb{N} \cup (4\pi \mathbb{N} + 4\pi(1+\alpha)).$$

**Consequence** (compactness): if  $\rho \notin \Lambda$ , then we have that

(1) solutions of  $(\tilde{E}_{\rho})$  stay bounded (and have bounded energy)

(2) the deformation lemma holds true (difference in topology of some sublevels implies existence of critical points of  $I_{\rho}$ )

$$\Lambda := 4\pi \mathbb{N} \cup (4\pi \mathbb{N} + 4\pi(1+\alpha)).$$

**Consequence** (compactness): if  $\rho \notin \Lambda$ , then we have that

(1) solutions of  $(\tilde{E}_{\rho})$  stay bounded (and have bounded energy)

(2) the deformation lemma holds true (difference in topology of some sublevels implies existence of critical points of  $I_{\rho}$ )

• Therefore, for  $\rho \notin \Lambda$ , it is crucial to understand the structure of sublevels of  $I_{\rho}$ 

$$\Lambda := 4\pi \mathbb{N} \cup (4\pi \mathbb{N} + 4\pi(1+\alpha)).$$

**Consequence** (compactness): if  $\rho \notin \Lambda$ , then we have that

(1) solutions of  $(\tilde{E}_{\rho})$  stay bounded (and have bounded energy)

(2) the deformation lemma holds true (difference in topology of some sublevels implies existence of critical points of  $I_{\rho}$ )

• Therefore, for  $\rho \notin \Lambda$ , it is crucial to understand the structure of sublevels of  $I_{\rho}$ 

For the regular case the above compactness theorem has previous counterparts by Y.Li, Shafrir, Ohtsuka, Suzuki, Ricciardi, C.Chen, C.S. Lin, ...

**Theorem 1** ([M.-Ruiz, '10]) Suppose  $\alpha \in (0,1]$  and that  $\rho \in (4\pi, 4\pi(1+\alpha))$ .

**Theorem 1** ([M.-Ruiz, '10]) Suppose  $\alpha \in (0,1]$  and that  $\rho \in (4\pi, 4\pi(1+\alpha))$ . Then, if  $\Sigma \not\simeq S^2$ , problem  $(E_{\rho})$  is solvable.

**Theorem 1** ([M.-Ruiz, '10]) Suppose  $\alpha \in (0,1]$  and that  $\rho \in (4\pi, 4\pi(1+\alpha))$ . Then, if  $\Sigma \not\simeq S^2$ , problem  $(E_{\rho})$  is solvable.

**Theorem 2** ([M.-Ruiz, '10]) Suppose  $\alpha \in (0,1)$  and that  $\rho \in (4\pi(1+\alpha), 8\pi)$ .

**Theorem 1** ([M.-Ruiz, '10]) Suppose  $\alpha \in (0,1]$  and that  $\rho \in (4\pi, 4\pi(1+\alpha))$ . Then, if  $\Sigma \not\simeq S^2$ , problem  $(E_{\rho})$  is solvable.

**Theorem 2** ([M.-Ruiz, '10]) Suppose  $\alpha \in (0,1)$  and that  $\rho \in (4\pi(1+\alpha), 8\pi)$ . Then problem  $(E_{\rho})$  is solvable for every  $\Sigma$ .

**Theorem 1** ([M.-Ruiz, '10]) Suppose  $\alpha \in (0,1]$  and that  $\rho \in (4\pi, 4\pi(1+\alpha))$ . Then, if  $\Sigma \not\simeq S^2$ , problem  $(E_{\rho})$  is solvable.

**Theorem 2** ([M.-Ruiz, '10]) Suppose  $\alpha \in (0,1)$  and that  $\rho \in (4\pi(1+\alpha), 8\pi)$ . Then problem  $(E_{\rho})$  is solvable for every  $\Sigma$ .

**Remark** ([Bartolucci - Lin - Tarantello, '10]) If  $(\Sigma, g) = (S^2, g_0)$ and if  $\rho \in (4\pi, 4\pi(1 + \alpha))$  then  $(E_{\rho})$  has no solution

**Theorem 1** ([M.-Ruiz, '10]) Suppose  $\alpha \in (0,1]$  and that  $\rho \in (4\pi, 4\pi(1+\alpha))$ . Then, if  $\Sigma \not\simeq S^2$ , problem  $(E_{\rho})$  is solvable.

**Theorem 2** ([M.-Ruiz, '10]) Suppose  $\alpha \in (0,1)$  and that  $\rho \in (4\pi(1+\alpha), 8\pi)$ . Then problem  $(E_{\rho})$  is solvable for every  $\Sigma$ .

**Remark** ([Bartolucci - Lin - Tarantello, '10]) If  $(\Sigma, g) = (S^2, g_0)$ and if  $\rho \in (4\pi, 4\pi(1 + \alpha))$  then  $(E_{\rho})$  has no solution, therefore our assumptions are somehow sharp.
## Some existence theorems (simple cases)

**Theorem 1** ([M.-Ruiz, '10]) Suppose  $\alpha \in (0,1]$  and that  $\rho \in (4\pi, 4\pi(1+\alpha))$ . Then, if  $\Sigma \not\simeq S^2$ , problem  $(E_{\rho})$  is solvable.

**Theorem 2** ([M.-Ruiz, '10]) Suppose  $\alpha \in (0,1)$  and that  $\rho \in (4\pi(1+\alpha), 8\pi)$ . Then problem  $(E_{\rho})$  is solvable for every  $\Sigma$ .

**Remark** ([Bartolucci - Lin - Tarantello, '10]) If  $(\Sigma, g) = (S^2, g_0)$ and if  $\rho \in (4\pi, 4\pi(1 + \alpha))$  then  $(E_{\rho})$  has no solution, therefore our assumptions are somehow sharp. The proof uses a Pohozaev type identity. Let  $L \in \mathbb{R}$  be sufficiently large:





 $\Rightarrow \qquad \{I_{\rho} \leq L\} \text{ has trivial topology} \qquad (L \gg 0)$ 



 $\Rightarrow \qquad \{I_{\rho} \leq L\} \text{ has trivial topology} \qquad (L \gg 0)$ 

**Corollary** If for some large L the *low sublevel*  $\{I_{\rho} \leq -L\}$  has non trivial topology, a solution to  $(\tilde{E}_{\rho})$  exists.

The Moser-Trudinger inequality can be improved for functions whose *conformal volume* is distributed into different regions.

The Moser-Trudinger inequality can be improved for functions whose *conformal volume* is distributed into different regions.

**Lemma 1** ([W.Chen - C. Li, '91]) Let  $S_1, S_2$  be subsets of  $\Sigma$  satisfying  $dist(S_1, S_2) \ge \delta_0$ , and let  $\gamma_0 > 0$ .

The Moser-Trudinger inequality can be improved for functions whose *conformal volume* is distributed into different regions.

**Lemma 1** ([W.Chen - C. Li, '91]) Let  $S_1, S_2$  be subsets of  $\Sigma$  satisfying  $dist(S_1, S_2) \ge \delta_0$ , and let  $\gamma_0 > 0$ . Suppose that



The Moser-Trudinger inequality can be improved for functions whose *conformal volume* is distributed into different regions.

**Lemma 1** ([W.Chen - C. Li, '91]) Let  $S_1, S_2$  be subsets of  $\Sigma$  satisfying  $dist(S_1, S_2) \ge \delta_0$ , and let  $\gamma_0 > 0$ . Suppose that



Then  $\forall \tilde{\varepsilon} > 0 \quad \exists$  a constant  $C = C(\tilde{\varepsilon}, \delta_0, \gamma_0)$  (indep. of u) s.t.

$$\log \int_{\Sigma} e^{2(u-\overline{u})} dV_g \le C + \frac{1}{8\pi - \tilde{\varepsilon}} \int_{\Sigma} |\nabla u|^2 dV_g$$

spreading  $\Rightarrow$  better const. in (M-T)  $\Rightarrow$  coercivity of  $I_{\rho}$ 

spreading  $\Rightarrow$  better const. in (M-T)  $\Rightarrow$  coercivity of  $I_{\rho}$  $I_{\rho}$  low  $\Rightarrow$  no better const. in (M-T)  $\Rightarrow$  no spreading

spreading  $\Rightarrow$  better const. in (M-T)  $\Rightarrow$  coercivity of  $I_{\rho}$ 

 $I_{\rho}$  low  $\Rightarrow$  no better const. in (M-T)  $\Rightarrow$  no spreading

**Lemma 2** For any  $\varepsilon > 0$  and any r > 0 there exists a large positive  $L = L(\varepsilon, r)$  such that for every  $u \in H^1(\Sigma)$  with  $I_{\rho}(u) \leq -L$  there exists a point  $p_u \in \Sigma$  such that  $\int_{\Sigma \setminus B_r(p_u)} h(x) e^{2u} dV_g < \varepsilon$ .

spreading  $\Rightarrow$  better const. in (M-T)  $\Rightarrow$  coercivity of  $I_{\rho}$ 

 $I_{\rho}$  low  $\Rightarrow$  no better const. in (M-T)  $\Rightarrow$  no spreading

**Lemma 2** For any  $\varepsilon > 0$  and any r > 0 there exists a large positive  $L = L(\varepsilon, r)$  such that for every  $u \in H^1(\Sigma)$  with  $I_{\rho}(u) \leq -L$  there exists a point  $p_u \in \Sigma$  such that  $\int_{\Sigma \setminus B_r(p_u)} h(x) e^{2u} dV_g < \varepsilon$ .



Therefore, for  $\rho < 8\pi$ , we obtain a natural continuous map

 $\Psi:\{I_{\rho}\leq -L\}\to \Sigma$ 

Therefore, for  $\rho < 8\pi$ , we obtain a natural continuous map

$$\Psi: \{I_{\rho} \leq -L\} \to \Sigma$$

which associates to any u with low energy the corresponding concentration point.

Therefore, for  $\rho < 8\pi$ , we obtain a natural continuous map

$$\Psi: \{I_{\rho} \leq -L\} \to \Sigma$$

which associates to any u with low energy the corresponding concentration point.

• We want to understand next what is the role of the singularity.

**Theorem** ([Dolbeault - Esteban - Tarantello, '08]) Suppose u is a radial function of class  $H_0^1(D)$ , where D is the unit disk of  $\mathbb{R}^2$ , and let  $\alpha > -1$ .

**Theorem** ([Dolbeault - Esteban - Tarantello, '08]) Suppose u is a radial function of class  $H_0^1(D)$ , where D is the unit disk of  $\mathbb{R}^2$ , and let  $\alpha > -1$ . Then there exists a constant  $C_{\alpha}$  such that

$$\int_D |x|^{2\alpha} e^{2u} dx \leq \frac{1}{4\pi(1+\alpha)} \int_D |\nabla u|^2 dx + C_\alpha.$$

**Theorem** ([Dolbeault - Esteban - Tarantello, '08]) Suppose u is a radial function of class  $H_0^1(D)$ , where D is the unit disk of  $\mathbb{R}^2$ , and let  $\alpha > -1$ . Then there exists a constant  $C_{\alpha}$  such that

$$\int_D |x|^{2\alpha} e^{2u} dx \leq \frac{1}{4\pi(1+\alpha)} \int_D |\nabla u|^2 dx + C_\alpha.$$

• Note that for  $\alpha > 0$  we now have a better constant in front of the Dirichlet norm.

**Theorem** ([Dolbeault - Esteban - Tarantello, '08]) Suppose u is a radial function of class  $H_0^1(D)$ , where D is the unit disk of  $\mathbb{R}^2$ , and let  $\alpha > -1$ . Then there exists a constant  $C_{\alpha}$  such that

$$\int_D |x|^{2\alpha} e^{2u} dx \leq \frac{1}{4\pi(1+\alpha)} \int_D |\nabla u|^2 dx + C_\alpha.$$

• Note that for  $\alpha > 0$  we now have a better constant in front of the Dirichlet norm. The proof uses a change of variable and the standard (M-T) inequality.

**Theorem** ([Dolbeault - Esteban - Tarantello, '08]) Suppose u is a radial function of class  $H_0^1(D)$ , where D is the unit disk of  $\mathbb{R}^2$ , and let  $\alpha > -1$ . Then there exists a constant  $C_{\alpha}$  such that

$$\int_D |x|^{2\alpha} e^{2u} dx \leq \frac{1}{4\pi(1+\alpha)} \int_D |\nabla u|^2 dx + C_\alpha.$$

• Note that for  $\alpha > 0$  we now have a better constant in front of the Dirichlet norm. The proof uses a change of variable and the standard (M-T) inequality. The constant is sharp, as one can see using modified  $\alpha$ -bubbles (giving the AF)

$$\varphi_{\alpha,\lambda}(x) = \log \frac{\lambda^{1+\alpha}}{1+(\lambda|x|)^{2(1+\alpha)}}$$

**Theorem** ([Dolbeault - Esteban - Tarantello, '08]) Suppose u is a radial function of class  $H_0^1(D)$ , where D is the unit disk of  $\mathbb{R}^2$ , and let  $\alpha > -1$ . Then there exists a constant  $C_{\alpha}$  such that

$$\int_D |x|^{2\alpha} e^{2u} dx \leq \frac{1}{4\pi(1+\alpha)} \int_D |\nabla u|^2 dx + C_\alpha.$$

• Note that for  $\alpha > 0$  we now have a better constant in front of the Dirichlet norm. The proof uses a change of variable and the standard (M-T) inequality. The constant is sharp, as one can see using modified  $\alpha$ -bubbles (giving the AF)

$$\varphi_{\alpha,\lambda}(x) = \log \frac{\lambda^{1+\alpha}}{1+(\lambda|x|)^{2(1+\alpha)}}$$

• The coefficient  $\frac{1}{4\pi(1+\alpha)}$  is the one to aim for, but radiality assumption is rather restrictive.

**Theorem** ([Dolbeault - Esteban - Tarantello, '08]) Suppose u is a radial function of class  $H_0^1(D)$ , where D is the unit disk of  $\mathbb{R}^2$ , and let  $\alpha > -1$ . Then there exists a constant  $C_{\alpha}$  such that

$$\int_D |x|^{2\alpha} e^{2u} dx \leq \frac{1}{4\pi(1+\alpha)} \int_D |\nabla u|^2 dx + C_\alpha.$$

• Note that for  $\alpha > 0$  we now have a better constant in front of the Dirichlet norm. The proof uses a change of variable and the standard (M-T) inequality. The constant is sharp, as one can see using modified  $\alpha$ -bubbles (giving the AF)

$$\varphi_{\alpha,\lambda}(x) = \log \frac{\lambda^{1+\alpha}}{1+(\lambda|x|)^{2(1+\alpha)}}.$$

• The coefficient  $\frac{1}{4\pi(1+\alpha)}$  is the one to aim for, but radiality assumption is rather restrictive. Our goal is to substitute it with a two-dimensional constraint, which is <u>much</u> more flexible.

A new Improved Moser-Trudinger inequality

**Proposition** Suppose  $\alpha \in (0, 1]$ .

**Proposition** Suppose  $\alpha \in (0, 1]$ . Then there exists L > 0 and a (continuous) *barycentric* map  $\beta : \{I_{\rho} \leq -L\} \rightarrow \Sigma$  such that

## A new Improved Moser-Trudinger inequality

**Proposition** Suppose  $\alpha \in (0, 1]$ . Then there exists L > 0 and a (continuous) *barycentric* map  $\beta : \{I_{\rho} \leq -L\} \rightarrow \Sigma$  such that for any  $\varepsilon > 0$ 

$$\log \int_{\Sigma} h(x)e^{2u} \leq \frac{1}{4\pi(1+\alpha)-\varepsilon} \int_{\Sigma} |\nabla u|^2 + 2 \int_{\Sigma} u + C_{\varepsilon},$$

### A new Improved Moser-Trudinger inequality

**Proposition** Suppose  $\alpha \in (0, 1]$ . Then there exists L > 0 and a (continuous) *barycentric* map  $\beta : \{I_{\rho} \leq -L\} \rightarrow \Sigma$  such that for any  $\varepsilon > 0$ 

$$\log \int_{\Sigma} h(x)e^{2u} \leq \frac{1}{4\pi(1+\alpha)-\varepsilon} \int_{\Sigma} |\nabla u|^2 + 2 \int_{\Sigma} u + C_{\varepsilon},$$
  
provided  $u \in \{I_{\rho} \leq -L\}$  and  $\beta(u) = p.$ 

**Proposition** Suppose  $\alpha \in (0, 1]$ . Then there exists L > 0 and a (continuous) *barycentric* map  $\beta : \{I_{\rho} \leq -L\} \rightarrow \Sigma$  such that for any  $\varepsilon > 0$ 

$$\log \int_{\Sigma} h(x) e^{2u} \leq \frac{1}{4\pi(1+\alpha) - \varepsilon} \int_{\Sigma} |\nabla u|^2 + 2 \int_{\Sigma} u + C_{\varepsilon},$$
  
provided  $u \in \{I_{\rho} \leq -L\}$  and  $\beta(u) = p.$ 

Idea of the proof.

### A new Improved Moser-Trudinger inequality

**Proposition** Suppose  $\alpha \in (0, 1]$ . Then there exists L > 0 and a (continuous) *barycentric* map  $\beta : \{I_{\rho} \leq -L\} \rightarrow \Sigma$  such that for any  $\varepsilon > 0$ 

$$\log \int_{\Sigma} h(x)e^{2u} \leq \frac{1}{4\pi(1+\alpha)-\varepsilon} \int_{\Sigma} |\nabla u|^2 + 2 \int_{\Sigma} u + C_{\varepsilon},$$
  
provided  $u \in \{I_{\rho} \leq -L\}$  and  $\beta(u) = p.$ 

Idea of the proof. First of all, set  $f = \frac{he^{2u}}{\int he^{2u}}$ .

**Proposition** Suppose  $\alpha \in (0, 1]$ . Then there exists L > 0 and a (continuous) *barycentric* map  $\beta : \{I_{\rho} \leq -L\} \rightarrow \Sigma$  such that for any  $\varepsilon > 0$ 

$$\log \int_{\Sigma} h(x)e^{2u} \leq \frac{1}{4\pi(1+\alpha)-\varepsilon} \int_{\Sigma} |\nabla u|^2 + 2 \int_{\Sigma} u + C_{\varepsilon},$$
  
provided  $u \in \{I_{\rho} \leq -L\}$  and  $\beta(u) = p.$ 

Idea of the proof. First of all, set  $f = \frac{he^{2u}}{\int he^{2u}}$ .

We next define a scale of concentration  $\sigma_x$  at every point x in the following way.
**Proposition** Suppose  $\alpha \in (0, 1]$ . Then there exists L > 0 and a (continuous) *barycentric* map  $\beta : \{I_{\rho} \leq -L\} \rightarrow \Sigma$  such that for any  $\varepsilon > 0$ 

$$\log \int_{\Sigma} h(x)e^{2u} \leq \frac{1}{4\pi(1+\alpha)-\varepsilon} \int_{\Sigma} |\nabla u|^2 + 2 \int_{\Sigma} u + C_{\varepsilon},$$
  
provided  $u \in \{I_{\rho} \leq -L\}$  and  $\beta(u) = p.$ 

Idea of the proof. First of all, set  $f = \frac{he^{2u}}{\int he^{2u}}$ .

We next define a scale of concentration  $\sigma_x$  at every point x in the following way. Fixing a large C > 0 let  $\sigma_x$  be such that

$$\int_{B_x(\sigma_x)} f = \int_{B_x(C\sigma_x)^c} f.$$

Let us now define

$$T(x) = \int_{B_x(\sigma_x)} f.$$

Let us now define

$$T(x) = \int_{B_x(\sigma_x)} f.$$

• Notice that both  $\sigma_x$  and T(x) are continuous (in x and u).

Let us now define

$$T(x) = \int_{B_x(\sigma_x)} f.$$

• Notice that both  $\sigma_x$  and T(x) are continuous (in x and u).

**Lemma** There exists  $\tau = \tau(C) > 0$  such that

$$\max_{x\in\Sigma} T(x) > 2\tau.$$

Let us now define

$$T(x) = \int_{B_x(\sigma_x)} f.$$

• Notice that both  $\sigma_x$  and T(x) are continuous (in x and u).

**Lemma** There exists  $\tau = \tau(C) > 0$  such that

$$\max_{x\in\Sigma} T(x) > 2\tau.$$

To see this, we use a covering argument

Let us now define

$$T(x) = \int_{B_x(\sigma_x)} f.$$

• Notice that both  $\sigma_x$  and T(x) are continuous (in x and u).

**Lemma** There exists  $\tau = \tau(C) > 0$  such that

$$\max_{x\in\Sigma} T(x) > 2\tau.$$

To see this, we use a covering argument (intuitively, there cannot be vanishing at all scales and at all points).

Let us now define

$$T(x) = \int_{B_x(\sigma_x)} f.$$

• Notice that both  $\sigma_x$  and T(x) are continuous (in x and u).

**Lemma** There exists  $\tau = \tau(C) > 0$  such that

$$\max_{x \in \Sigma} T(x) > 2\tau.$$

To see this, we use a covering argument (intuitively, there cannot be vanishing at all scales and at all points).

Define then:

$$\beta: \{I_{\rho} \leq -L\} \rightarrow \Sigma; \qquad \beta(u) = \frac{\int_{\Sigma} [T(x,f) - \tau]^{+} x}{\int_{\Sigma} [T(x,f) - \tau]^{+}}.$$

Recall that

$$\beta(u) = \frac{\int_{\Sigma} [T(x,f) - \tau]^+ x}{\int_{\Sigma} [T(x,f) - \tau]^+}.$$

Recall that 
$$\beta(u) = \frac{\int_{\Sigma} [T(x, f) - \tau]^+ x}{\int_{\Sigma} [T(x, f) - \tau]^+}.$$

Suppose that  $\beta(u) = p$ , and define  $\overline{\sigma}$  as

$$\overline{\sigma} = \sup \left\{ \sigma_x : T(x) \ge \tau \right\}.$$

Recall that 
$$\beta(u) = \frac{\int_{\Sigma} [T(x, f) - \tau]^+ x}{\int_{\Sigma} [T(x, f) - \tau]^+}.$$

Suppose that  $\beta(u) = p$ , and define  $\overline{\sigma}$  as

$$\overline{\sigma} = \sup \left\{ \sigma_x : T(x) \ge \tau \right\}.$$

Recall that 
$$\beta(u) = \frac{\int_{\Sigma} [T(x, f) - \tau]^+ x}{\int_{\Sigma} [T(x, f) - \tau]^+}.$$

Suppose that  $\beta(u) = p$ , and define  $\overline{\sigma}$  as

$$\overline{\sigma} = \sup \left\{ \sigma_x : T(x) \ge \tau \right\}.$$

**1)** 
$$\int_{B_y(\overline{\sigma})} f \geq \tau;$$

Recall that 
$$\beta(u) = \frac{\int_{\Sigma} [T(x,f) - \tau]^+ x}{\int_{\Sigma} [T(x,f) - \tau]^+}.$$

Suppose that  $\beta(u) = p$ , and define  $\overline{\sigma}$  as

$$\overline{\sigma} = \sup \left\{ \sigma_x : T(x) \ge \tau \right\}.$$

**1)** 
$$\int_{B_y(\overline{\sigma})} f \ge \tau$$
; **2)**  $\int_{B_y(C\overline{\sigma})^c} f \ge \tau$ ;

Recall that 
$$\beta(u) = \frac{\int_{\Sigma} [T(x, f) - \tau]^+ x}{\int_{\Sigma} [T(x, f) - \tau]^+}.$$

Suppose that  $\beta(u) = p$ , and define  $\overline{\sigma}$  as

$$\overline{\sigma} = \sup \left\{ \sigma_x : T(x) \ge \tau \right\}.$$

**1)** 
$$\int_{B_y(\overline{\sigma})} f \ge \tau$$
; **2)**  $\int_{B_y(C\overline{\sigma})^c} f \ge \tau$ ; **3)**  $dist(y,p) \le C\overline{\sigma}$ 

Recall that 
$$\beta(u) = \frac{\int_{\Sigma} [T(x, f) - \tau]^+ x}{\int_{\Sigma} [T(x, f) - \tau]^+}$$

٠

Suppose that  $\beta(u) = p$ , and define  $\overline{\sigma}$  as

$$\overline{\sigma} = \sup \left\{ \sigma_x : T(x) \ge \tau \right\}.$$

Then there exists  $y \in \Sigma$  such that

**1)** 
$$\int_{B_y(\overline{\sigma})} f \ge \tau$$
; **2)**  $\int_{B_y(C\overline{\sigma})^c} f \ge \tau$ ; **3)**  $dist(y,p) \le C\overline{\sigma}$ 

**Good news:** we now have concentration at scale  $\overline{\sigma}$  and at distance  $\overline{\sigma}$  from the singularity p

Recall that 
$$\beta(u) = \frac{\int_{\Sigma} [T(x, f) - \tau]^+ x}{\int_{\Sigma} [T(x, f) - \tau]^+}$$

Suppose that  $\beta(u) = p$ , and define  $\overline{\sigma}$  as

$$\overline{\sigma} = \sup \left\{ \sigma_x : T(x) \ge \tau \right\}.$$

Then there exists  $y \in \Sigma$  such that

**1)** 
$$\int_{B_y(\overline{\sigma})} f \ge \tau$$
; **2)**  $\int_{B_y(C\overline{\sigma})^c} f \ge \tau$ ; **3)**  $dist(y,p) \le C\overline{\sigma}$ 

**Good news:** we now have concentration at scale  $\overline{\sigma}$  and at distance  $\overline{\sigma}$  from the singularity p (notice that  $\overline{\sigma}$  might not be continuous in u).

With 1)-3) we can now estimate the log of the integral in two ways.

With 1)-3) we can now estimate the log of the integral in two ways. The first uses the inner inequality and a scaling

$$\log \int_{\Sigma} h(x) e^{2u} \leq C_{\tau} + \log \int_{B_y(\overline{\sigma})} |x|^{2\alpha} e^{2u}$$

With 1)-3) we can now estimate the log of the integral in two ways. The first uses the inner inequality and a scaling

$$\log \int_{\Sigma} h(x) e^{2u} \leq C_{\tau} + \log \int_{B_{y}(\overline{\sigma})} |x|^{2\alpha} e^{2u} \leq C_{\tau} + \frac{1}{4\pi} \int_{B_{y}(\overline{\sigma})} |\nabla u|^{2} + 2\alpha \log \overline{\sigma} + \int_{\partial B_{p}(\overline{\sigma})} u.$$

With 1)-3) we can now estimate the log of the integral in two ways. The first uses the inner inequality and a scaling

$$\log \int_{\Sigma} h(x) e^{2u} \leq C_{\tau} + \log \int_{B_y(\overline{\sigma})} |x|^{2\alpha} e^{2u} \leq C_{\tau} + \frac{1}{4\pi} \int_{B_y(\overline{\sigma})} |\nabla u|^2 + 2\alpha \log \overline{\sigma} + \int_{\partial B_p(\overline{\sigma})} u.$$

The second uses the outer inequality

$$\log \int_{\Sigma} h(x) e^{2u} \leq C_{ au} + \log \int_{B_y(C\overline{\sigma})^c} |x|^{2lpha} e^{2u}.$$

With 1)-3) we can now estimate the log of the integral in two ways. The first uses the inner inequality and a scaling

$$\log \int_{\Sigma} h(x) e^{2u} \leq C_{\tau} + \log \int_{B_{y}(\overline{\sigma})} |x|^{2\alpha} e^{2u} \leq C_{\tau} + \frac{1}{4\pi} \int_{B_{y}(\overline{\sigma})} |\nabla u|^{2} + 2\alpha \log \overline{\sigma} + \int_{\partial B_{p}(\overline{\sigma})} u.$$

The second uses the outer inequality

$$\log \int_{\Sigma} h(x) e^{2u} \leq C_{ au} + \log \int_{B_y(C\overline{\sigma})^c} |x|^{2lpha} e^{2u}.$$

Setting  $v = u + 2\alpha \log |x|$ , we estimate the last integral applying (M-T) to v

With 1)-3) we can now estimate the log of the integral in two ways. The first uses the inner inequality and a scaling

$$\log \int_{\Sigma} h(x) e^{2u} \leq C_{\tau} + \log \int_{B_{y}(\overline{\sigma})} |x|^{2\alpha} e^{2u} \leq C_{\tau} + \frac{1}{4\pi} \int_{B_{y}(\overline{\sigma})} |\nabla u|^{2} + 2\alpha \log \overline{\sigma} + \int_{\partial B_{p}(\overline{\sigma})} u.$$

The second uses the outer inequality

$$\log \int_{\Sigma} h(x) e^{2u} \leq C_{ au} + \log \int_{B_y(C\overline{\sigma})^c} |x|^{2lpha} e^{2u}.$$

Setting  $v = u + 2\alpha \log |x|$ , we estimate the last integral applying (M-T) to v

$$\log \int_{B_y(C\overline{\sigma})^c} |x|^{2lpha} e^{2u} \leq -2\log \overline{\sigma} + rac{1}{4\pi} \int_{B_y(C\overline{\sigma})^c} |
abla v|^2.$$

With 1)-3) we can now estimate the log of the integral in two ways. The first uses the inner inequality and a scaling

$$\log \int_{\Sigma} h(x) e^{2u} \leq C_{\tau} + \log \int_{B_{y}(\overline{\sigma})} |x|^{2\alpha} e^{2u} \leq C_{\tau} + \frac{1}{4\pi} \int_{B_{y}(\overline{\sigma})} |\nabla u|^{2} + 2\alpha \log \overline{\sigma} + \int_{\partial B_{p}(\overline{\sigma})} u.$$

The second uses the outer inequality

$$\log \int_{\Sigma} h(x) e^{2u} \leq C_{ au} + \log \int_{B_y(C\overline{\sigma})^c} |x|^{2lpha} e^{2u}.$$

Setting  $v = u + 2\alpha \log |x|$ , we estimate the last integral applying (M-T) to v

$$\log \int_{B_y(C\overline{\sigma})^c} |x|^{2\alpha} e^{2u} \leq -2\log \overline{\sigma} + \frac{1}{4\pi} \int_{B_y(C\overline{\sigma})^c} |\nabla v|^2.$$

The terms  $\nabla u \cdot \nabla \log |x|$  and  $|\nabla (\log |x|)|^2$  give back the boundary integral  $\int_{\partial B_p(\overline{\sigma})} u$  and  $\log \overline{\sigma}$  with the right coefficients.

With 1)-3) we can now estimate the log of the integral in two ways. The first uses the inner inequality and a scaling

$$\log \int_{\Sigma} h(x) e^{2u} \leq C_{\tau} + \log \int_{B_{y}(\overline{\sigma})} |x|^{2\alpha} e^{2u} \leq C_{\tau} + \frac{1}{4\pi} \int_{B_{y}(\overline{\sigma})} |\nabla u|^{2} + 2\alpha \log \overline{\sigma} + \int_{\partial B_{p}(\overline{\sigma})} u.$$

The second uses the outer inequality

$$\log \int_{\Sigma} h(x) e^{2u} \leq C_{ au} + \log \int_{B_y(C\overline{\sigma})^c} |x|^{2lpha} e^{2u}.$$

Setting  $v = u + 2\alpha \log |x|$ , we estimate the last integral applying (M-T) to v

$$\log \int_{B_y(C\overline{\sigma})^c} |x|^{2lpha} e^{2u} \leq -2\log \overline{\sigma} + rac{1}{4\pi} \int_{B_y(C\overline{\sigma})^c} |
abla v|^2.$$

The terms  $\nabla u \cdot \nabla \log |x|$  and  $|\nabla (\log |x|)|^2$  give back the boundary integral  $\int_{\partial B_p(\overline{\sigma})} u$  and  $\log \overline{\sigma}$  with the right coefficients. **q.e.d.** 

• Choosing a scale and a point of *maximal concentration* by looking at suitable integrals is a standard procedure in geometric analysis.

• Choosing a scale and a point of maximal concentration by looking at suitable integrals is a standard procedure in geometric analysis. The issue here is to obtain <u>continuity</u> of the barycentric map  $\beta$ , which is fundamental for us.

• Choosing a scale and a point of maximal concentration by looking at suitable integrals is a standard procedure in geometric analysis. The issue here is to obtain <u>continuity</u> of the barycentric map  $\beta$ , which is fundamental for us.

• The new feature of our inequality is that it is scaling invariant.

• Choosing a scale and a point of maximal concentration by looking at suitable integrals is a standard procedure in geometric analysis. The issue here is to obtain <u>continuity</u> of the barycentric map  $\beta$ , which is fundamental for us.

• The new feature of our inequality is that it is <u>scaling invariant</u>. In the previous improvements (to our knowledge) a lower bound on distances was always needed.

• Choosing a scale and a point of maximal concentration by looking at suitable integrals is a standard procedure in geometric analysis. The issue here is to obtain <u>continuity</u> of the barycentric map  $\beta$ , which is fundamental for us.

• The new feature of our inequality is that it is <u>scaling invariant</u>. In the previous improvements (to our knowledge) a lower bound on distances was always needed.

• The assumption  $\alpha \leq 1$  is sharp.

• Choosing a scale and a point of maximal concentration by looking at suitable integrals is a standard procedure in geometric analysis. The issue here is to obtain <u>continuity</u> of the barycentric map  $\beta$ , which is fundamental for us.

• The new feature of our inequality is that it is <u>scaling invariant</u>. In the previous improvements (to our knowledge) a lower bound on distances was always needed.

• The assumption  $\alpha \leq 1$  is sharp. For  $\alpha > 1$  *splitting* the mass gives a worse constant and lowers the functional.

## Proof of Theorem 1: $\rho \in (4\pi, 4\pi(1 + \alpha))$

# Proof of Theorem 1: $\rho \in (4\pi, 4\pi(1 + \alpha))$

For  $\delta > 0$  small consider the family of functions

$$\varphi_{\lambda,x}(y) := \log \frac{\lambda}{1 + \lambda^2 dist(x,y)^2}; \qquad x \in \Sigma \setminus B_{\delta}(p).$$

## Proof of Theorem 1: $\rho \in (4\pi, 4\pi(1 + \alpha))$

For  $\delta > 0$  small consider the family of functions

$$\varphi_{\lambda,x}(y) := \log \frac{\lambda}{1 + \lambda^2 dist(x,y)^2}; \qquad x \in \Sigma \setminus B_{\delta}(p).$$

Then

 $I_{\rho}(\varphi_{\lambda,x}) \to -\infty \text{ as } \lambda \to +\infty \quad \text{ uniformly for } x \in \Sigma \setminus B_{\delta}(p)$
For  $\delta > 0$  small consider the family of functions

$$\varphi_{\lambda,x}(y) := \log \frac{\lambda}{1 + \lambda^2 dist(x,y)^2}; \qquad x \in \Sigma \setminus B_{\delta}(p).$$

Then

 $I_{\rho}(\varphi_{\lambda,x}) \to -\infty$  as  $\lambda \to +\infty$  uniformly for  $x \in \Sigma \setminus B_{\delta}(p)$ , so we obtain a map from  $\Sigma \setminus B_{\delta}(p)$  into low sublevels of  $I_{\rho}$ .

For  $\delta > 0$  small consider the family of functions

$$\varphi_{\lambda,x}(y) := \log \frac{\lambda}{1 + \lambda^2 dist(x,y)^2}; \qquad x \in \Sigma \setminus B_{\delta}(p).$$

Then

 $I_{\rho}(\varphi_{\lambda,x}) \to -\infty$  as  $\lambda \to +\infty$  uniformly for  $x \in \Sigma \setminus B_{\delta}(p)$ , so we obtain a map from  $\Sigma \setminus B_{\delta}(p)$  into low sublevels of  $I_{\rho}$ .

Viceversa, the barycentric map  $\beta$  on low sublevels cannot be equal to p, by the improved inequality.

For  $\delta > 0$  small consider the family of functions

$$\varphi_{\lambda,x}(y) := \log \frac{\lambda}{1 + \lambda^2 dist(x,y)^2}; \qquad x \in \Sigma \setminus B_{\delta}(p).$$

Then

 $I_{\rho}(\varphi_{\lambda,x}) \to -\infty$  as  $\lambda \to +\infty$  uniformly for  $x \in \Sigma \setminus B_{\delta}(p)$ , so we obtain a map from  $\Sigma \setminus B_{\delta}(p)$  into low sublevels of  $I_{\rho}$ .

Viceversa, the barycentric map  $\beta$  on low sublevels cannot be equal to p, by the improved inequality.

Therefore we obtain a reverse map from low sublevels of  $I_{\rho}$  into  $\Sigma \setminus \{p\} \simeq \Sigma \setminus B_{\delta}(p)$ .

For  $\delta > 0$  small consider the family of functions

$$\varphi_{\lambda,x}(y) := \log \frac{\lambda}{1 + \lambda^2 dist(x,y)^2}; \qquad x \in \Sigma \setminus B_{\delta}(p).$$

Then

 $I_{\rho}(\varphi_{\lambda,x}) \to -\infty$  as  $\lambda \to +\infty$  uniformly for  $x \in \Sigma \setminus B_{\delta}(p)$ , so we obtain a map from  $\Sigma \setminus B_{\delta}(p)$  into low sublevels of  $I_{\rho}$ .

Viceversa, the barycentric map  $\beta$  on low sublevels cannot be equal to p, by the improved inequality.

Therefore we obtain a reverse map from low sublevels of  $I_{\rho}$  into  $\Sigma \setminus \{p\} \simeq \Sigma \setminus B_{\delta}(p)$ .

When  $\Sigma$  is not the sphere  $\Sigma \setminus B_{\delta}(p)$  is not contractible, so we obtain that  $\{I_{\rho} \leq -L\}$ , L large, has non trivial topology.

For  $\delta > 0$  small consider the family of functions

$$\varphi_{\lambda,x}(y) := \log \frac{\lambda}{1 + \lambda^2 dist(x,y)^2}; \qquad x \in \Sigma \setminus B_{\delta}(p).$$

Then

 $I_{\rho}(\varphi_{\lambda,x}) \to -\infty$  as  $\lambda \to +\infty$  uniformly for  $x \in \Sigma \setminus B_{\delta}(p)$ , so we obtain a map from  $\Sigma \setminus B_{\delta}(p)$  into low sublevels of  $I_{\rho}$ .

Viceversa, the barycentric map  $\beta$  on low sublevels cannot be equal to p, by the improved inequality.

Therefore we obtain a reverse map from low sublevels of  $I_{\rho}$  into  $\Sigma \setminus \{p\} \simeq \Sigma \setminus B_{\delta}(p)$ .

When  $\Sigma$  is not the sphere  $\Sigma \setminus B_{\delta}(p)$  is not contractible, so we obtain that  $\{I_{\rho} \leq -L\}$ , L large, has non trivial topology. **q.e.d.** 

It is sufficient to modify the previous argument including the point p and using the modified test functions

$$\varphi_{\alpha,\lambda,x}(y) := \log \frac{\lambda^{1+\alpha}}{1+\lambda^2 dist(x,y)^{2(1+\alpha)}}; \qquad x \in \Sigma.$$

It is sufficient to modify the previous argument including the point p and using the modified test functions

$$\varphi_{\alpha,\lambda,x}(y) := \log \frac{\lambda^{1+\alpha}}{1+\lambda^2 dist(x,y)^{2(1+\alpha)}}; \qquad x \in \Sigma.$$

Now one has

 $I_{\rho}(\varphi_{\alpha,\lambda,x}) \to -\infty \text{ as } \lambda \to +\infty \quad \text{ uniformly for } \underline{\text{every }} x \in \Sigma.$ 

It is sufficient to modify the previous argument including the point p and using the modified test functions

$$\varphi_{\alpha,\lambda,x}(y) := \log \frac{\lambda^{1+\alpha}}{1+\lambda^2 dist(x,y)^{2(1+\alpha)}}; \qquad x \in \Sigma$$

Now one has

 $I_{\rho}(\varphi_{\alpha,\lambda,x}) \to -\infty \text{ as } \lambda \to +\infty \quad \text{ uniformly for } \underline{\text{every }} x \in \Sigma.$ 

The conclusion follows from the fact that every compact surface is non contractible.

It is sufficient to modify the previous argument including the point p and using the modified test functions

$$\varphi_{\alpha,\lambda,x}(y) := \log \frac{\lambda^{1+\alpha}}{1+\lambda^2 dist(x,y)^{2(1+\alpha)}}; \qquad x \in \Sigma$$

Now one has

 $I_{\rho}(\varphi_{\alpha,\lambda,x}) \to -\infty \text{ as } \lambda \to +\infty \quad \text{ uniformly for } \underline{\text{every }} x \in \Sigma.$ 

The conclusion follows from the fact that every compact surface is non contractible. **q.e.d.** 

For more singularities, when  $\rho \notin \Lambda$  (whose definition has to be suitably modified), there are still ways to prove existence when

For more singularities, when  $\rho \notin \Lambda$  (whose definition has to be suitably modified), there are still ways to prove existence when **a)**  $\alpha_i > 0$  for all *i* and  $\Sigma$  is not the sphere.

For more singularities, when  $\rho \notin \Lambda$  (whose definition has to be suitably modified), there are still ways to prove existence when

a)  $\alpha_i > 0$  for all *i* and  $\Sigma$  is not the sphere. In [Bartolucci-De Marchis-M, '10] using only partial topological information.

For more singularities, when  $\rho \notin \Lambda$  (whose definition has to be suitably modified), there are still ways to prove existence when

a)  $\alpha_i > 0$  for all *i* and  $\Sigma$  is not the sphere. In [Bartolucci-De Marchis-M, '10] using only partial topological information.

**b)**  $\alpha_i < 0$  for all *i* under further restrictions ([Carlotto-M, '11]).

For more singularities, when  $\rho \notin \Lambda$  (whose definition has to be suitably modified), there are still ways to prove existence when

**a)**  $\alpha_i > 0$  for all *i* and  $\Sigma$  is not the sphere. In [Bartolucci-De Marchis-M, '10] using only partial topological information.

**b)**  $\alpha_i < 0$  for all *i* under further restrictions ([Carlotto-M, '11]).

• The real issue is to get an improved inequality for  $\alpha > 1$ , which should be substantially different.

For more singularities, when  $\rho \notin \Lambda$  (whose definition has to be suitably modified), there are still ways to prove existence when

a)  $\alpha_i > 0$  for all *i* and  $\Sigma$  is not the sphere. In [Bartolucci-De Marchis-M, '10] using only partial topological information.

**b)**  $\alpha_i < 0$  for all *i* under further restrictions ([Carlotto-M, '11]).

• The real issue is to get an improved inequality for  $\alpha > 1$ , which should be substantially different.

• In many geometrically relevant situations (prescription of the Gaussian curvature)  $\rho$  might belong to the critical set  $\Lambda$ , by the Gauss-Bonnet formula.

For more singularities, when  $\rho \notin \Lambda$  (whose definition has to be suitably modified), there are still ways to prove existence when

**a)**  $\alpha_i > 0$  for all *i* and  $\Sigma$  is not the sphere. In [Bartolucci-De Marchis-M, '10] using only partial topological information.

**b)**  $\alpha_i < 0$  for all *i* under further restrictions ([Carlotto-M, '11]).

• The real issue is to get an improved inequality for  $\alpha > 1$ , which should be substantially different.

• In many geometrically relevant situations (prescription of the Gaussian curvature)  $\rho$  might belong to the critical set  $\Lambda$ , by the Gauss-Bonnet formula.

In this case one would need a more refined blow-up analysis, together with some detailed information on the Green's function or some linear combination of different ones.

The study of the selfdual Chern-Simons equations in the non abelian case leads to coupled Toda systems

The study of the selfdual Chern-Simons equations in the non abelian case leads to coupled Toda systems, of the form

$$\begin{cases} -\Delta u_1 + 2\rho_1 - \rho_2 = 2\rho_1 h_1(x)e^{u_1} - \rho_2 h_2(x)e^{u_2}; \\ -\Delta u_2 + 2\rho_2 - \rho_1 = 2\rho_2 h_2(x)e^{u_2} - \rho_1 h_1(x)e^{u_1} \end{cases}$$

The study of the selfdual Chern-Simons equations in the non abelian case leads to coupled Toda systems, of the form

$$\begin{cases} -\Delta u_1 + 2\rho_1 - \rho_2 = 2\rho_1 h_1(x)e^{u_1} - \rho_2 h_2(x)e^{u_2}; \\ -\Delta u_2 + 2\rho_2 - \rho_1 = 2\rho_2 h_2(x)e^{u_2} - \rho_1 h_1(x)e^{u_1}, \end{cases}$$

possibly with singularities on the right-hand side.

The study of the selfdual Chern-Simons equations in the non abelian case leads to coupled Toda systems, of the form

$$\begin{cases} -\Delta u_1 + 2\rho_1 - \rho_2 = 2\rho_1 h_1(x)e^{u_1} - \rho_2 h_2(x)e^{u_2}; \\ -\Delta u_2 + 2\rho_2 - \rho_1 = 2\rho_2 h_2(x)e^{u_2} - \rho_1 h_1(x)e^{u_1}, \end{cases}$$

possibly with singularities on the right-hand side.

Some results are available concerning classification and blow-up analysis of solutions ([Jost-Wang], [Jost-Lin-Wang], ...)

The study of the selfdual Chern-Simons equations in the non abelian case leads to coupled Toda systems, of the form

$$\begin{cases} -\Delta u_1 + 2\rho_1 - \rho_2 = 2\rho_1 h_1(x)e^{u_1} - \rho_2 h_2(x)e^{u_2}; \\ -\Delta u_2 + 2\rho_2 - \rho_1 = 2\rho_2 h_2(x)e^{u_2} - \rho_1 h_1(x)e^{u_1}, \end{cases}$$

possibly with singularities on the right-hand side.

Some results are available concerning classification and blow-up analysis of solutions ([Jost-Wang], [Jost-Lin-Wang], ...), but the existence question is almost entirely open.

The study of the selfdual Chern-Simons equations in the non abelian case leads to coupled Toda systems, of the form

$$\begin{cases} -\Delta u_1 + 2\rho_1 - \rho_2 = 2\rho_1 h_1(x)e^{u_1} - \rho_2 h_2(x)e^{u_2}; \\ -\Delta u_2 + 2\rho_2 - \rho_1 = 2\rho_2 h_2(x)e^{u_2} - \rho_1 h_1(x)e^{u_1}, \end{cases}$$

possibly with singularities on the right-hand side.

Some results are available concerning classification and blow-up analysis of solutions ([Jost-Wang], [Jost-Lin-Wang], ...), but the existence question is almost entirely open. From many aspects the singular scalar equation seems like a *toy model* for understanding this system.

The study of the selfdual Chern-Simons equations in the non abelian case leads to coupled Toda systems, of the form

$$\begin{cases} -\Delta u_1 + 2\rho_1 - \rho_2 = 2\rho_1 h_1(x)e^{u_1} - \rho_2 h_2(x)e^{u_2}; \\ -\Delta u_2 + 2\rho_2 - \rho_1 = 2\rho_2 h_2(x)e^{u_2} - \rho_1 h_1(x)e^{u_1}, \end{cases}$$

possibly with singularities on the right-hand side.

Some results are available concerning classification and blow-up analysis of solutions ([Jost-Wang], [Jost-Lin-Wang], ...), but the existence question is almost entirely open. From many aspects the singular scalar equation seems like a *toy model* for understanding this system.

• The problem without the self-duality assumption seems for the moment completely out of reach.

Thank you for your attention